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Optical Induction in Chiral Crystals. II.* The Crystal and Molecular Structures of 1-(2,6-Dichlorophenyl)-4-phenyl-*trans,trans*-1,3-butadiene

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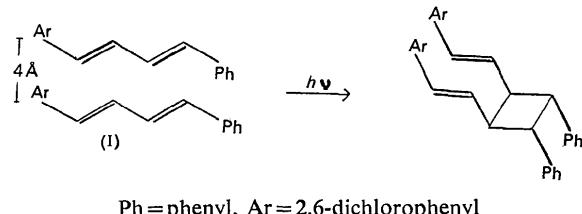
The title compound ($C_{16}H_{12}Cl_2$) crystallizes in space group $P2_12_12_1$ with $a=9.3881$ (2), $b=4.0286$ (1) and $c=35.3073$ (3) Å, $Z=4$. The structure has been determined by direct methods from three-dimensional diffractometer data collected with $Cu K\alpha$ radiation, and refined to $R=0.04$. The absolute configuration of the molecules in the analysed crystal was determined by the anomalous scattering of the chlorine atoms. The rotations of the phenyl and the dichlorophenyl groups from the butadiene plane are 15° and 40° respectively. The molecular conformation is believed to control the exclusive photoreactivity of the double bond adjacent to the phenyl ring and the formation of an optically active heterodimer in a mixed single crystal of the title compound with its thiophene analogue.

Introduction

Two of the necessary conditions for optical synthesis discussed in part I (Rabinovich & Shakked, 1974), chiral space group and *prochiral* conformation, were observed in the designing of another asymmetric synthesis experiment, namely, the photodimerization of parallel C=C double bonds stacked along a short (*ca* 4 Å) axis. Irradiation of such a system yields a cyclobutane possessing a mirror plane (Schmidt, 1967). However, as was suggested by Lahav (1968), the irradiation of a chiral single crystal containing a guest molecular species may lead to the formation of unequal amounts of enantiomeric heterodimers if the molecules are heavily distorted from planarity.

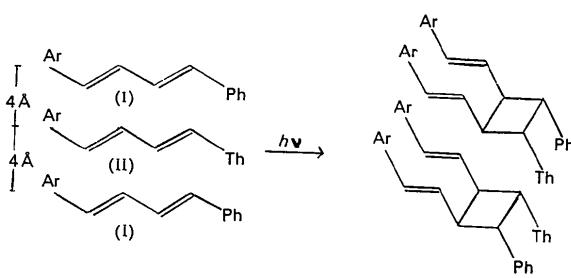
The present analysis deals with the crystal and molecular structures of 1-(2,6-dichlorophenyl)-4-phenyl-*trans,trans*-1,3-butadiene (I) which crystallizes in a chiral space group ($P2_12_12_1$) with a short (4.03 Å) axis and is expected to possess a pronounced *prochiral* con-

formation due to the two *ortho*-substituted chlorines. This compound yields under irradiation a single cyclobutane resulting from exclusive dimerization at the double bond adjacent to the phenyl ring (Cohen, Elgavi, Green, Ludmer & Schmidt, 1972).



(I) and its thiophene analogue [1-(2,6-dichlorophenyl)-4-thienyl-*trans,trans*-1,3-butadiene, (II)] are isomorphous and show extensive solid solubility. Irradiation of a mixed single crystal of (I) and (II) (approximate ratio 4:1) in the thiophene absorbing wavelength range yields one of the two enantiomeric heterodimers in excess (Elgavi, 1973; Elgavi, Green & Schmidt, 1973).

* Part I: Rabinovich & Shakked (1974).



Ph=phenyl, Th=2-thienyl, Ar=2,6-dichlorophenyl

The structure analysis of (I) was needed in order to explain both the exclusive photoreactivity of one of the two available double bonds and the formation of an optically active product in the mixed crystal system.

Experimental

Crystals of (I) (Fig. 1) were grown by slow evaporation of ethyl acetate solutions. The colourless crystals are elongated along [010]. A nearly equidimensional crystal was obtained by cleaving a needle-shaped crystal.

The crystal was mounted along [010] on a General Electric diffractometer. Accurate cell dimensions were derived by a least-squares procedure using high-order reflexions measured with Cu $K\alpha$ radiation ($2\theta > 120^\circ$). Crystal data are listed in Table 1.

Table 1. *Crystal data*

1-(2,6-Dichlorophenyl)-4-phenyl-trans,trans-1,3-butadiene,
 $C_{16}H_{12}Cl_2$, m.p. 108°
Orthorhombic, $P2_12_12_1$ M.W. 275.1
 $a = 9.3881 (2) \text{ \AA}$ $D_x = 1.37 \text{ g cm}^{-3}$ for $Z = 4$
 $b = 4.0286 (2)$ $D_m = 1.36 \text{ g cm}^{-3}$
 $c = 35.3073 (3)$ $\mu(\text{Cu } K\alpha) = 41.6 \text{ cm}^{-1}$

Three-dimensional intensity data ($2\theta \leq 160^\circ$, one octant of the reciprocal sphere, 1787 reflexions) were

collected manually by the stationary-counter stationary-crystal technique. The measurements were made at room temperature with Ni-filtered Cu $K\alpha$ radiation. The intensities of four standard reflexions were monitored every 100 reflexions throughout the data collection. The overall reduction in the intensity due to crystal deterioration was about 6%. The intensities were corrected for Lorentz, polarization and absorption effects. The weighting factors $w(hkl)$ and assignment of unobserved reflexions were derived by a procedure previously described (Filippakis, Leiserowitz & Schmidt, 1967).

Structure solution and refinement

The structure was solved by a multi-solution technique utilizing the tangent formula (Kaufman & Leiserowitz, 1970). An E map based on the phases of the model with the lowest (0.29) R_{Karle} (Karle & Karle, 1966) yielded all the non-hydrogen atoms in the molecule. These atoms were refined isotropically by a full-matrix least-squares program. Eventually, all hydrogens were inserted in chemically reasonable positions and the structure was further refined (anisotropically

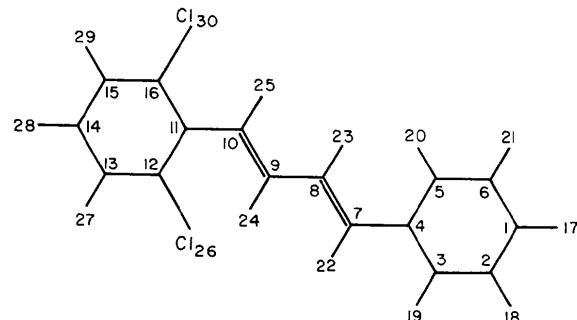


Fig. 1. 1-(2,6-Dichlorophenyl)-4-phenyl-trans,trans-1,3-butadiene (I). Chemical formula and numbering of atoms.

Table 2. *Fractional coordinates and vibration tensor components (\AA^2) of the heavy atoms and their e.s.d.'s*

The anisotropic temperature factor is of the form: $\exp[-2\pi^2(h_i a^i)(h_j a^j) U^{ij} \times 10^{-4}]$, where h_i and a^i , $i=1,2,3$ are reflexion indices and reciprocal unit-cell edges respectively.

	$x \times 10^4$	$y \times 10^4$	$z \times 10^5$	U^{11}	U^{22}	U^{33}	U^{12}	U^{13}	U^{23}
C(1)	2218 (4)	3647 (11)	-8966 (8)	726 (23)	711 (27)	394 (15)	-76 (25)	-39 (19)	12 (16)
C(2)	1017 (4)	5311 (10)	-7750 (8)	604 (20)	698 (27)	462 (16)	-25 (22)	86 (19)	-81 (15)
C(3)	849 (3)	6077 (9)	-3965 (8)	524 (18)	565 (24)	483 (16)	29 (20)	31 (17)	-29 (14)
C(4)	1865 (3)	5142 (9)	-1297 (7)	473 (16)	466 (20)	368 (13)	-53 (16)	26 (15)	-10 (12)
C(5)	3053 (3)	3402 (10)	-2570 (8)	462 (16)	577 (23)	466 (15)	13 (19)	3 (17)	-53 (13)
C(6)	3233 (4)	2683 (11)	-6360 (9)	557 (20)	730 (29)	507 (17)	5 (21)	-116 (20)	55 (15)
C(7)	1634 (3)	6113 (10)	2665 (8)	512 (17)	519 (23)	437 (15)	-8 (19)	-32 (16)	38 (13)
C(8)	2332 (3)	4971 (9)	5717 (8)	554 (18)	495 (19)	378 (14)	-82 (19)	-12 (15)	22 (13)
C(9)	2015 (3)	6212 (9)	9490 (8)	586 (18)	493 (22)	405 (14)	-50 (19)	-9 (16)	13 (14)
C(10)	2653 (3)	5135 (10)	12624 (8)	608 (20)	514 (21)	400 (15)	-38 (20)	-14 (16)	40 (14)
C(11)	2441 (3)	6465 (9)	16468 (7)	609 (17)	433 (18)	364 (13)	-48 (18)	29 (14)	-20 (12)
C(12)	1123 (3)	7413 (9)	17988 (8)	611 (19)	558 (23)	408 (15)	-52 (19)	-1 (16)	-24 (14)
C(13)	981 (4)	8755 (11)	21582 (9)	682 (21)	658 (26)	492 (17)	-9 (24)	-19 (20)	51 (17)
C(14)	2168 (4)	9140 (11)	23829 (9)	888 (26)	752 (31)	436 (17)	-10 (26)	-11 (20)	17 (18)
C(15)	3483 (4)	8141 (12)	22533 (9)	830 (26)	791 (29)	407 (17)	-54 (27)	-30 (20)	-10 (17)
C(16)	3592 (3)	6848 (10)	18928 (8)	615 (20)	571 (23)	445 (15)	29 (20)	60 (18)	-26 (14)
Cl(26)	-444 (1)	6817 (3)	15496 (2)	590 (5)	1183 (9)	623 (5)	-79 (7)	-147 (6)	-8 (4)
Cl(30)	5290 (1)	5703 (3)	17434 (2)	642 (5)	1193 (10)	531 (4)	148 (7)	29 (6)	-55 (4)

for carbon and chlorine, and isotropically for the hydrogen atoms) by a block-diagonal least-squares program to the following agreement values:

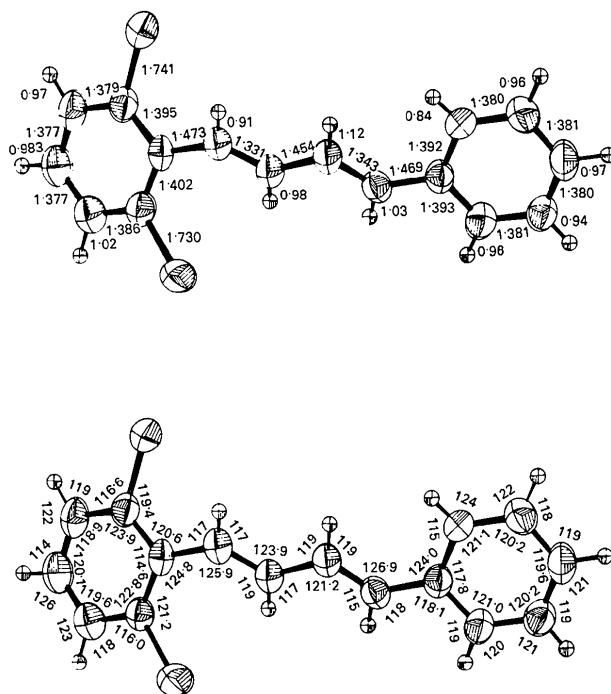


Fig. 2. Bond lengths (\AA) and angles ($^\circ$). The thermal ellipsoids are drawn at 50% probability level (Johnson, 1965). The view is along the b axis.

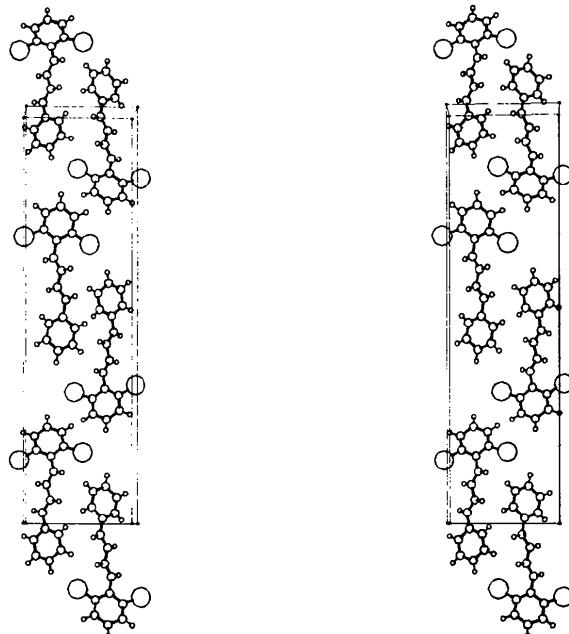


Fig. 3. Stereoscopic view of the packing arrangement along the b axis. The origin of the unit cell is in the bottom left-hand corner with c axis vertical and a axis horizontal and b entering the paper.

$$R = \sum kF_o - |F_c| / \sum kF_o = 0.040$$

$$r = \sum w(k^2 F_o^2 - |F_c|^2)^2 / \sum wk^4 F_o^4 = 0.006.$$

The number of refined parameters and the number of observations included in the last cycle of refinement were 211 and 1520 respectively.

The scattering factors used are those of Berghuis, Haanappel, Potters, Loopstra, MacGillavry & Veenendaal (1955) for C, Dawson (1960) for Cl and McWeeny (1951) for H. Next, anomalous dispersion corrections were introduced for Cl, the values used ($\Delta f' = 0.3$, $\Delta f'' = 0.7$) being taken from *International Tables for X-ray Crystallography* (1962). Refinements of both enantiomorphic structures were carried on until the shifts of the parameters were less than 0.2σ . Final agreement values of the two enantiomorphs are as follows:

$$R = 0.037, 0.047 \\ r = 0.004, 0.008.$$

The coordinates and thermal parameters of the two enantiomeric structures agree within 2σ .

Use of the *R* test (Hamilton, 1965) indicates that the difference is highly significant and the choice of handedness implied by the parameters in Table 2 with a right-handed coordinate system is correct. In this context however (of flexible molecules undergoing spontaneous resolution) the chirality indicated by these results is merely a consequence of the arbitrary choice of a crystal; it has no chemical significance. The final parameters and their e.s.d.'s are listed in Tables 2 and 3. A list of observed and calculated structure factors is given in Table 4.

Table 3. Hydrogen atom parameters and their e.s.d.'s

	$x (\times 10^3)$	$y (\times 10^3)$	$z (\times 10^4)$	$\bar{U}^2 (\times 10^3 \text{ \AA}^2)$
H(17)	236 (3)	313 (9)	-1162 (8)	33 (9)
H(18)	33 (3)	598 (9)	-953 (8)	33 (10)
H(19)	3 (3)	726 (9)	-312 (8)	35 (11)
H(20)	365 (3)	295 (10)	-88 (7)	25 (10)
H(21)	408 (3)	161 (11)	-729 (9)	56 (12)
H(22)	92 (3)	780 (10)	319 (8)	34 (10)
H(23)	304 (3)	273 (8)	539 (7)	12 (8)
H(24)	130 (3)	797 (11)	970 (8)	39 (11)
H(25)	335 (3)	358 (9)	1232 (7)	24 (9)
H(27)	-1 (3)	942 (9)	2245 (7)	31 (10)
H(28)	218 (4)	1001 (11)	2643 (10)	67 (13)
H(29)	434 (3)	842 (9)	2404 (8)	37 (10)

Results and discussion

Thermal motion analysis

Least-squares analysis of the possible rigid-body motion of the molecule was carried out by the method of Schomaker & Trueblood (1968). Three rigid-body models were used, comprising different groups of atoms: the first model consisting of all non-hydrogen atoms, the second, of the atoms of the dichlorophenyl group and the central butadiene chain, and the third of the phenyl ring and the butadiene chain. The r.m.s.

values of ΔU^{ij} (corrected for the number of degrees of freedom) obtained for the three models are: 0.0053, 0.0049 and 0.0039 Å² respectively, significantly higher than the mean value (0.0018 Å²) of $\sigma(U^{ij})$. These figures show that none of the models properly represents rigid-body motion. The best fit is obtained for the third model, indicating that the internal vibrations centre mainly at the single bond adjacent to the dichlorophenyl group. The results of the analysis of the

first model are given in Table 5. The translational vibrations are nearly isotropic; the librational motion is highly anisotropic with the largest libration (5.5°) about the longest molecular axis. The results of the analysis of the second model are similar. The third model, however, yields a smaller librational motion (4.7°) about the longest axis of the molecule, thus substantiating our former indication about the internal vibrations of the molecule.

Table 4. Observed and calculated structure factors

Each set of three columns contains respectively l , $10F_o$ and $10F_c$. Unobserved reflexions are marked by U and zero-weighted reflexions by an asterisk.

Table 5. Results of rigid-body motion analysis

The eigenvectors of \mathbf{L} and \mathbf{T} are referred to the molecular system of inertia \mathbf{I} (in terms of the corresponding direction cosines $\times 10^4$). The eigenvectors of \mathbf{I} are referred to the reciprocal vectors a^* , b^* and c^* . Components of the origin shifts $\varrho(\varrho_1, \varrho_2, \varrho_3)$, referred to the molecular centroids are given in Å.

Eigenvalues		Eigenvectors			
663 (at. wt. Å ²)		0.601	0.533	34.923	
3721	I	9.251	-0.666	-1.470	
4337		1.483	3.937	-4.978	
(5.5°)*	92($\times 10^{-4}$ rad ²)	9997	-97	211	
(1.6)	8	L	96	-6536	-7568
(1.4)	6		212	7568	-6534
(0.22 Å)	502($\times 10^{-4}$ Å ²)	-1827	9314	-3150	
(0.21)	436	T	-893	-2911	-3429
(0.20)	403		-411	2186	8850

$$\varrho(-1.44, 0.01, 0.12) \\ [(\sum (\Delta U^{ij})^2 / (6n-s)]^{1/2} = 0.0053 \text{ Å}^2 (n=18, s=20)$$

* The corresponding r.m.s. amplitudes are given in parentheses.

Molecular geometry

Fig. 2 shows the experimental bond lengths and angles, uncorrected for librational motion. The mean values of the e.s.d.'s in the bond lengths of C-Cl, C-C and C-H are 0.003, 0.004 and 0.03 Å respectively. The corresponding values for bond angles are 0.2, 0.3 and 1.8°. The 'bond-scatter' $\{\sum(r_i - \bar{r})^2 / (n-1)\}^{1/2}$, $\bar{r} = 1.385 \text{ Å}$ of the 12 aromatic bonds (0.008 Å) is significantly larger than the average e.s.d. (0.004 Å), probably due to the relative elongation of the four aromatic bonds adjacent to the butadiene chain. The 'bond scatter' value obtained without these bonds is 0.003 Å. The C-Cl bond lengths (1.741, 1.730 Å) agree well with the mean value ($1.737 \pm 0.016 \text{ Å}$) given by Palenik, Donohue & Trueblood (1968).

The 'angle scatter' (defined analogously to the 'bond scatter') calculated from the 12 aromatic bond angles (2.3°) is much larger than the average e.s.d. (0.3°), indicating significant deviations from 120°. The largest deviations are observed at C(12), C(16) and C(11). The

first two values (122.8, 123.9°) are attributed to the chlorine substituents, since an aromatic substitution by electron-withdrawing groups generally leads to an increase in the internal angle at the substituted position. This effect has been observed in chlorine substitution as well (Hope, 1969; Bernstein, 1972). The remarkably small value (114.6°) of the aromatic angle at C(11) results from the substitution effect of the conjugated chain similarly to that observed at C(4) of the phenyl ring (117.8°), and from the narrowing which takes place to compensate for the widening at C(12) and C(16). Similar decreased values are observed in other *meta*-disubstituted benzene derivatives by electron-withdrawing groups (Trotter & Williston, 1966; Palenik *et al.*, 1968; Hope, 1969; Bavoux & Thozet, 1973).

The central C-C bond (1.454 Å) of the conjugated chain shows a decreased single-bond character, compared to the peripheral single bonds (1.469, 1.473 Å), in accordance with other molecules containing conjugated chains (e.g., Stam & Riva di Sanseverino, 1966; Filippakis *et al.*, 1967). The bond and torsion angles (Table 6) of the central chain are affected mainly by the aromatic rings which form sterically hindered conjugated systems. The relaxation of the steric interactions near the phenyl ring is achieved by the opening of the bond angles at C(4) and C(7), (124.0, 126.9°) and a twist of 14.4° about C(4)-C(7), as observed in some chalcone derivatives (Rabinovich & Shakked, 1974). The non-bonded interactions near the dichlorophenyl ring are relaxed mainly by similar increase in the bond angles (125.9, 124.8°) at C(10) and C(11) and a large twist (40.8°) about C(10)-C(11). The significantly larger value (123.9°) of C(8)-C(9)-C(10) compared to that (121.2°) of the equivalent angle C(7)-C(8)-C(9) results also from these interactions by 'chain stretching'.

The minimum-energy geometry of the molecule in the gaseous and crystalline phases was calculated by means of a computer program (Warshel, 1974) based on the quantum mechanical extension of the Con-

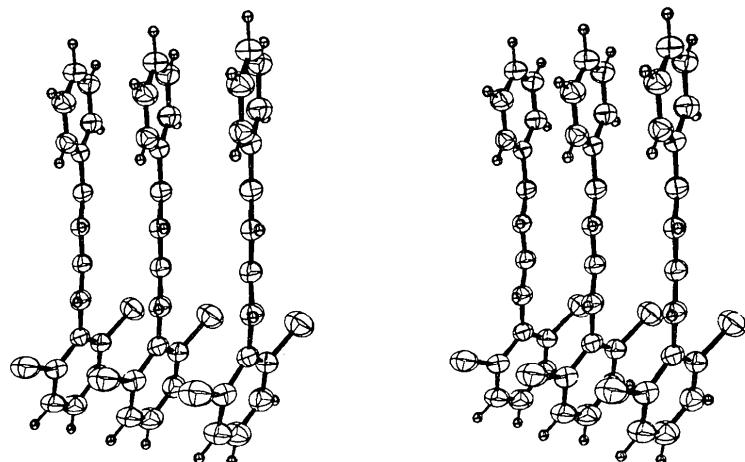


Fig. 4. Stereoscopic view of the molecules in the 4 Å translation stack. The view is edge-on along the plane of the butadiene.

Table 6. Selected bond lengths (\AA), bond angles ($^\circ$) and torsion angles ($^\circ$) of the calculated and observed molecular geometry

The calculations were carried out by the QCFF/PI scheme (Warshel, 1974). The observed values are corrected for libration. A positive torsion angle corresponds to a right-handed screw (Klyne & Prelog, 1960).

	Gas (calc.)	Crystal (calc.)	Crystal (obs.)
Bond			
C(3)—C(4)	1.420	1.419	1.397
C(4)—C(5)	1.418	1.417	1.398
C(4)—C(7)	1.480	1.478	1.470
C(7)—C(8)	1.355	1.354	1.346
C(8)—C(9)	1.469	1.468	1.456
C(9)—C(10)	1.351	1.351	1.334
C(10)—C(11)	1.487	1.486	1.474
C(11)—C(12)	1.423	1.424	1.409
C(11)—C(16)	1.422	1.422	1.399
Bond angle			
C(3)—C(4)—C(7)	118.5	118.5	118.1
C(5)—C(4)—C(7)	122.9	122.8	124.0
C(4)—C(7)—C(8)	125.3	125.2	126.9
C(7)—C(8)—C(9)	120.7	120.6	121.2
C(8)—C(9)—C(10)	122.3	122.3	123.9
C(9)—C(10)—C(11)	123.5	123.7	125.9
C(10)—C(11)—C(12)	122.2	122.6	124.8
C(10)—C(11)—C(16)	119.1	119.0	120.6
C(11)—C(12)—Cl(26)	122.8	123.1	121.2
C(11)—C(16)—Cl(30)	121.7	121.9	119.4
Torsion angle			
C(5)—C(4)—C(7)—C(8)	-0.4	-15.1	-14.4
C(4)—C(7)—C(8)—C(9)	179.3	177.9	178.6
C(7)—C(8)—C(9)—C(10)	181.4	179.0	178.8
C(8)—C(9)—C(10)—C(11)	177.8	176.0	175.6
C(9)—C(10)—C(11)—C(12)	43.5	39.7	40.8
C(10)—C(11)—C(12)—Cl(26)	5.0	6.6	4.6
C(10)—C(11)—C(16)—Cl(30)	-0.3	-1.3	-0.6

sistent Force-Field to π electron system (Warshel & Karplus, 1972). The energy parameters used, involving C and H atoms were those of set *A* given by Warshel, Huller, Rabinovich & Shakked (1974) and the non-bonded potential parameters of Cl \cdots Cl, Cl \cdots C and Cl \cdots H were the refined values of set *B* given by Bonadeo & D'Alessio (1973). Some of the calculated results are given in Table 6 together with the corresponding experimental values corrected for libration motion. The agreement between the observed and calculated values is fairly good, considering the approximations and assumptions introduced into the calculations on one hand, and the difficulty of correcting properly for librational motion, on the other. The relatively large deviations (2.5°) between the calculated and observed bond angles at C(11), C(12) and C(16) are due to the fact that no account was taken of the electron-withdrawing effects of the chlorine atoms in decreasing the *s*-character contribution to the related σ -bonds (Carter, McPhail & Sim, 1966). A comparison between the molecular geometry in the gaseous and the crystalline phase, leads to the conclusion that the twist about the single bond C(4)—C(7) results from crystal forces whereas the twist about C(10)—C(11) is due mainly to intramolecular forces.

Molecular packing

Fig. 3 shows a stereoscopic view of the packing arrangement along [010]. The molecules pack in stacks along the short axis. The short intermolecular contacts are in the range of the quoted van der Waals distances (see Table 7). The shortest Cl \cdots Cl contact (4.03 \AA) is between atoms related by translation along the *b* axis, and is significantly longer than the sum of the van der Waals radii (3.6 \AA). As stated above the 4 \AA packing mode is a necessary topochemical condition for photocycloaddition reaction. This packing mode can be achieved with high degree of probability in aromatic compounds varying in size, symmetry and molecular weight, by introducing two chlorine atoms into a phenyl ring as shown by Green, Leser & Schmidt (1973). Similar effects were also observed in dibromo and diiodo compounds and are attributed to the dispersion energy gained by the interaction of two pairs of halogen atoms at this separation.

Table 7. Short intermolecular distances (\AA)

The second atom in each pair is related to the first by the corresponding symmetry and translation operations.

	Symmetry element	Translation along		
		<i>a</i>	<i>b</i>	<i>c</i>
C(1) \cdots Cl(26)	2	0	1	0
C(3) \cdots H(23)	2	-1	0	0
C(4) \cdots C(5)	1	0	1	0
C(5) \cdots H(22)	2	0	1	0
C(8) \cdots H(19)	2	0	1	0
C(13) \cdots H(27)	3	0	-1	0
C(14) \cdots H(27)	3	0	-1	0
C(15) \cdots H(29)	3	1	-1	0
H(17) \cdots Cl(26)	2	0	0	0
H(17) \cdots Cl(26)	2	0	1	0
H(17) \cdots Cl(30)	2	-1	0	0
H(18) \cdots Cl(30)	2	-1	1	0
H(21) \cdots Cl(26)	2	0	0	0
H(28) \cdots Cl(30)	3	1	0	0
H(29) \cdots H(29)	3	1	0	0
H(29) \cdots Cl(30)	3	1	0	0

Symmetry elements: (1) x, y, z ; (2) $\frac{1}{2}+x, \frac{1}{2}-y, -z$; (3) $-x, \frac{1}{2}+y, \frac{1}{2}-z$.

Photoreactivity

Fig. 4 presents a stereoscopic view of the molecules in the 4 \AA translation stack, seen edge-on along the butadiene plane. This view demonstrates the steric effects of the dichlorophenyl groups which prevent close approach of adjacent molecules, and may explain the exclusive dimerization at the C=C double bond adjacent to the unsubstituted phenyl ring.

The contact geometry in the ground electronic state of any molecule with the adjacent molecules above it and below it along [010] is identical owing to translation symmetry. However, the excited molecule does not necessarily possess the same equilibrium geometry as in the ground state (Suzuki, 1967) and hence, any geometrical relaxation of such a non-planar molecule in the excited electronic state will lead to different ap-

proach configurations with its ground-state partner above and below, of which one may be more favourable for excimer formation (Cohen, Elgavi, Green, Ludmer & Schmidt, 1972) and subsequent dimerization. This behaviour is a necessary condition for creating unequal amounts of enantiomeric heterodimers in the two-component system (I and II) where the guest molecule (II), interleaved between two host molecules (I), is selectively excited as described earlier. An attempt is being made to study the preferred direction of photodimerization in this system by theoretical methods on one hand, and on the other by comparing the absolute configuration of the molecules in the initial mixed crystal with that of the heterodimer obtained in excess on irradiation of the same crystal.

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